

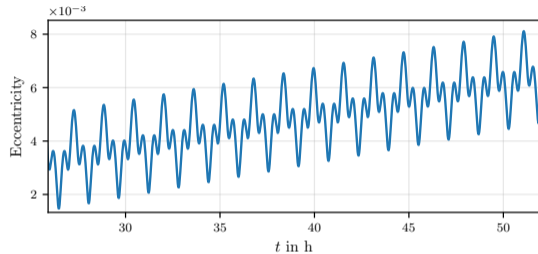
Efficient Numerical Optimal Control for Highly Oscillatory Systems

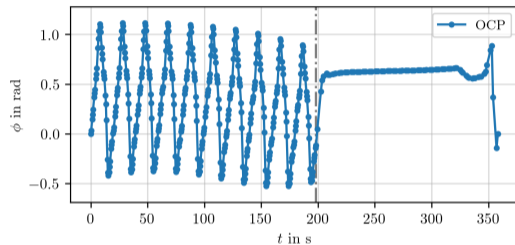
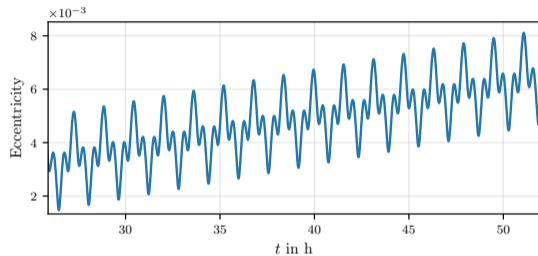
Jakob Harzer, Jochem De Schutter, Moritz Diehl

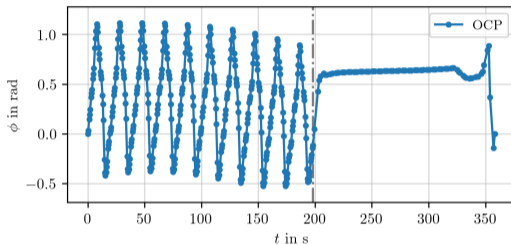
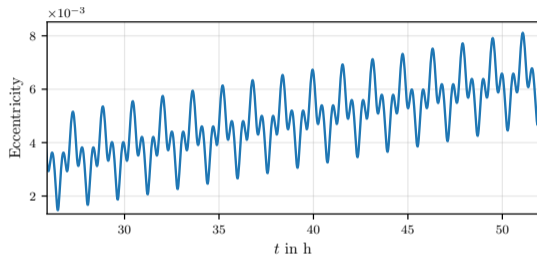
CDC 2022 Presentation
December 6th, 2022



Introduction







Can we make use of the fact that the expected trajectory of an OCP is highly oscillatory to speed up the optimization process?



- ▶ Highly oscillatory initial value problem

$$\dot{x}(t) = f(x(t))$$

$$x(0) = x_0$$

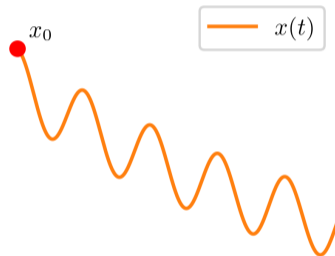
with solution $x(t) \in \mathbb{R}^{n_x}, t \in [0, t_f]$ and known period duration T .

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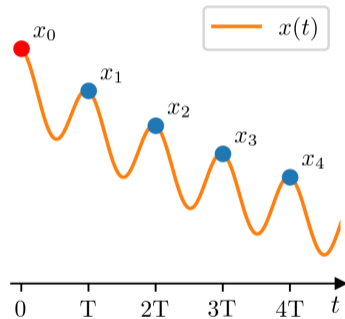


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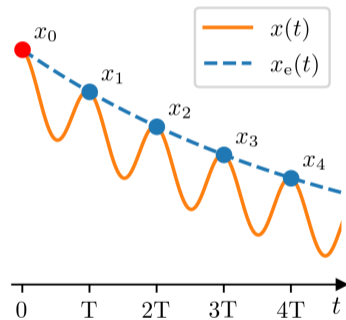
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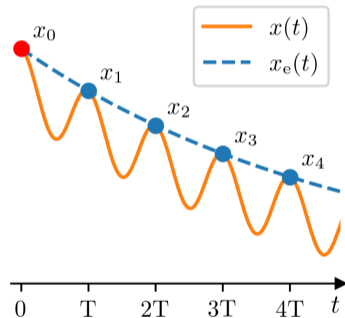
with solution $x(t) \in \mathbb{R}^{n_x}, t \in [0, t_f]$ and known period duration T .

- ▶ Construct envelope $x_e(t) \in \mathbb{R}^{n_x}$ that smoothly interpolates a series of 'periodic' points x_1, x_2, \dots



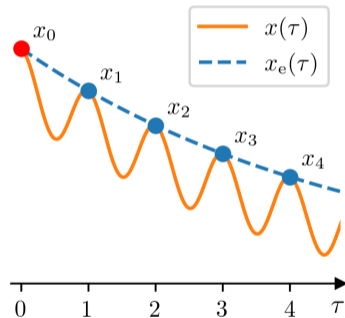
- ▶ Assumption: Envelope x_e stems from envelope dynamics

$$\frac{dx_e}{dt}(t) = f_e(x_e(t))$$



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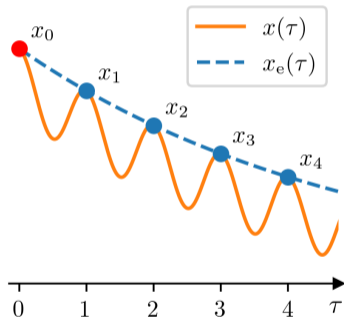
Envelope Dynamics



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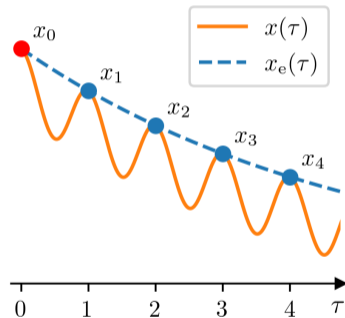


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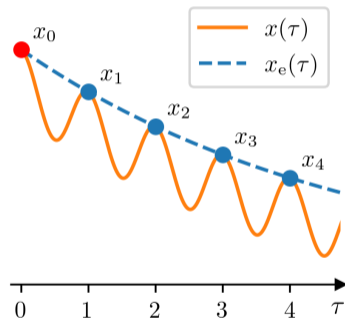
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1. Micro-integrate one period of the oscillations
2. Approximate the envelope dynamics using a DAE

$$\begin{aligned}\frac{dx_e}{d\tau}(\tau) &= f_e(x_e(\tau), z_e(\tau)) \\ 0 &= g_e(x_e(\tau), z_e(\tau))\end{aligned}$$



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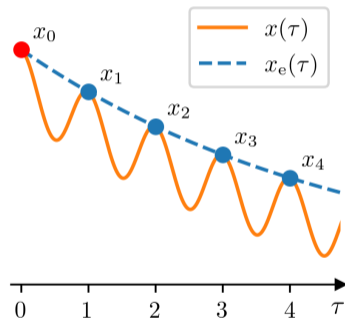
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3. Macro-integrate the approximated slow envelope dynamics instead of the fast oscillating dynamics!





"Basic underlying idea has appeared several times in the literature over the last fifty years" [3]

- ▶ Multirevolution Methods by Mace and Thomas[5], Graf[4]
- ▶ Envelope Following Methods by Petzold[6]
- ▶ Stroboscopic Averaging Methods by Calvo[2]
- ▶ Heterogeneous Multiscale Methods, Averaging Methods etc.



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New: Fully implicit DAE formulation, Use for Optimal Control Problems

Toy Example



Toy Example



- ▶ Perturbed Predator-Prey Model with rabbits r and snow leopards s :

$$x = \begin{bmatrix} r \\ s \end{bmatrix} \in \mathbb{R}^2$$

$$\dot{x} = \begin{bmatrix} \alpha r - \beta r s \\ \delta r s - \gamma s \end{bmatrix} + \begin{bmatrix} 0 \\ u \end{bmatrix}$$

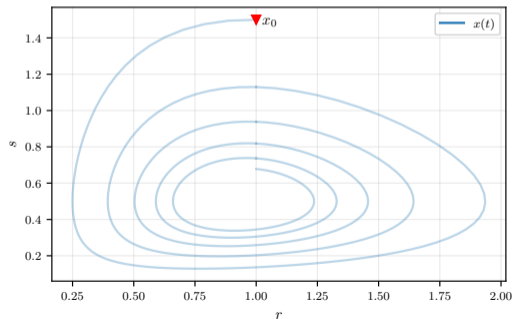
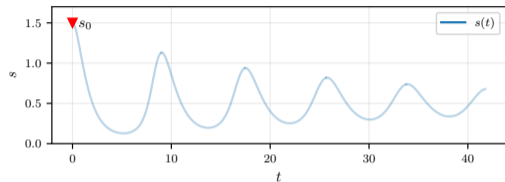
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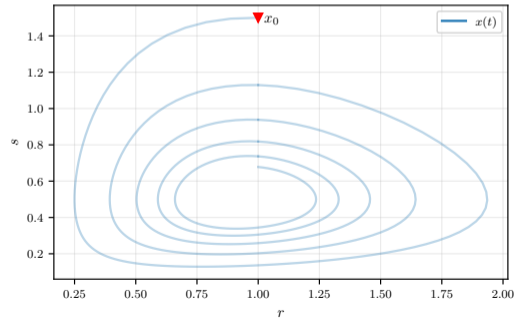
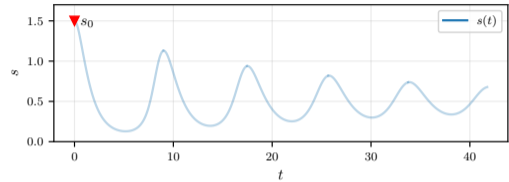
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Cycle Discretization



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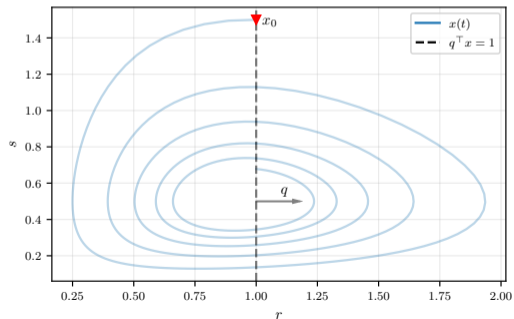
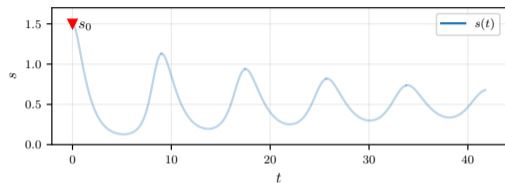


► Affine phase conditions

$$q^\top x^- = b^- \quad (1a)$$

$$q^\top x^+ = b^+ \quad (1b)$$

for the start point $x^- \in \mathbb{R}^{n_x}$ and end point $x^+ \in \mathbb{R}^{n_x}$ of a cycle, respectively.



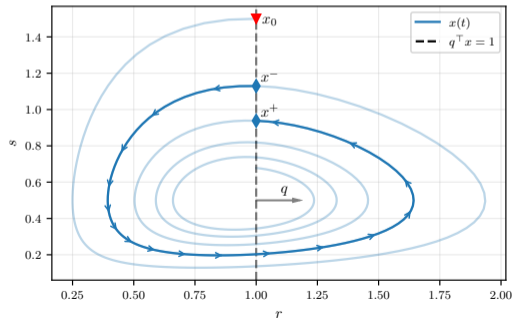
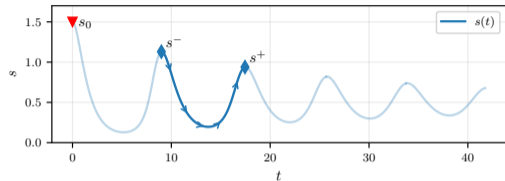
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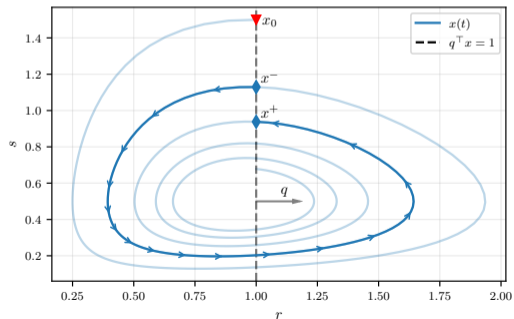
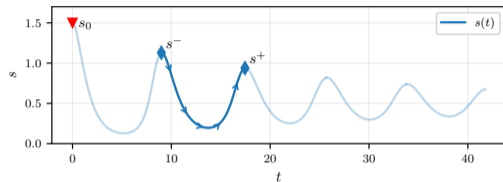
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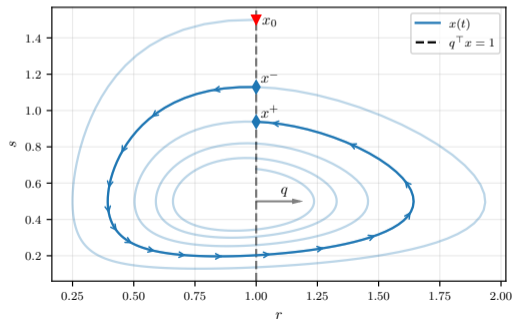
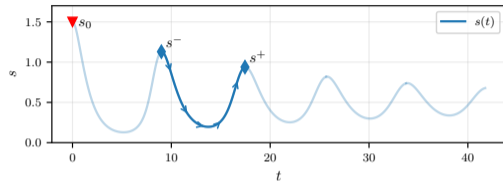
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- ▶ $[Q|q]$ is an orthonormal basis of the statespace



Cycle Discretization II



- ▶ Discrete start and end point of a cycle x_k^-, x_k^+



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$$C(z_k, U_k) := \begin{cases} 0 & = q^\top x_k^- - b^- \\ 0 & = x_k^+ - F(x_k^-, U_k, T_k) \\ 0 & = q^\top x_k^+ - b^+ \end{cases} \in \mathbb{R}^{n_x+2} \quad (2)$$

that define one cycle.

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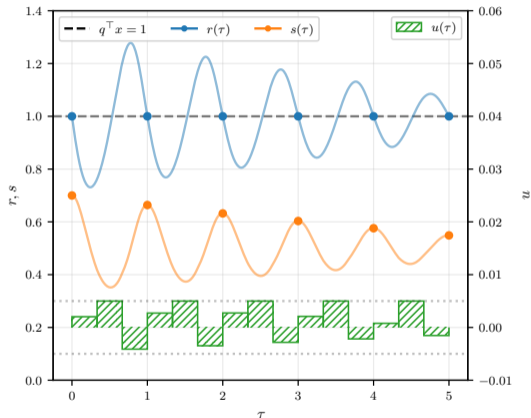
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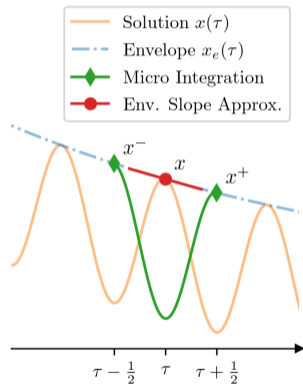
- ▶ Algebraic cycle variables

$$z_k = (x_k^-, x_k^+, T_k)$$

$$\begin{aligned}
 \min_w \quad & \sum_{k=0}^{N-1} L_c(z_k, U_k) + E(x_{N-1}^+) \\
 \text{s.t.} \quad & 0 = Q^\top(x_0^- - x_0), \\
 & 0 = C(z_k, U_k) \quad k = 0, \dots, N-1, \\
 & 0 = Q^\top(x_k^- - x_{k-1}^+) \quad k = 1, \dots, N-1
 \end{aligned}$$



Central Difference Envelope-DAE



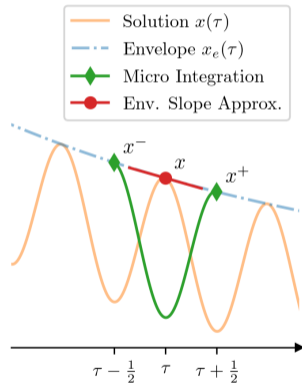


► Central Difference Envelope-DAE

$$\frac{dx}{d\tau} = x^+ - x^- \quad (\text{slope approx.})$$

$$0 = Q^\top \left(\frac{x^+ + x^-}{2} - x \right) \quad (\text{connecting cond.})$$

$$0 = C(z, U) \quad (\text{cycle cond.})$$



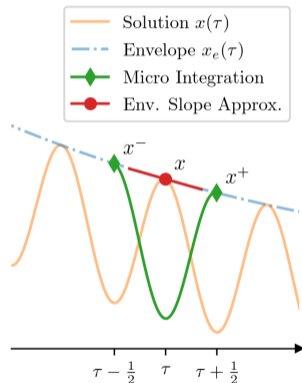
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- ▶ summarize by $f_e(x, z) := x^+ - x^-$ and $0 = g_e(x, z)$ with $x = x$ and $z = (x^+, x^-, T)$



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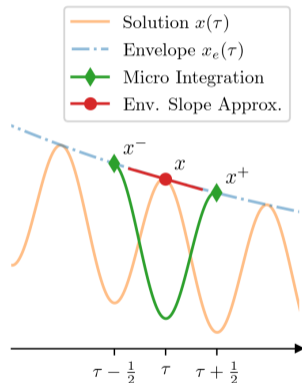
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▶ Macro-Integration of the Envelope-DAE over $\tau \in [0, N]$

$$\frac{dx}{d\tau}(\tau) = f_e(x(\tau), z(\tau))$$

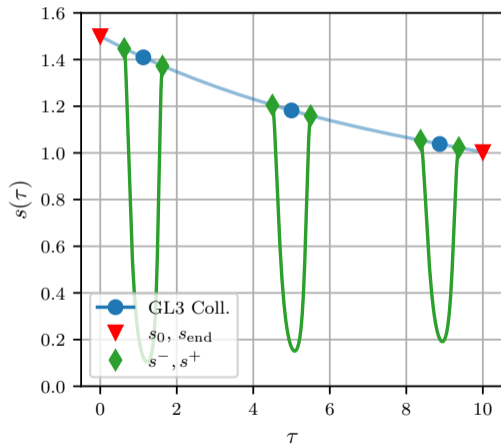
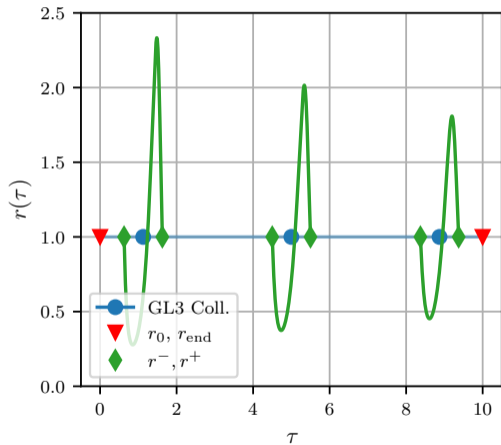
$$0 = g_e(x(\tau), z(\tau))$$



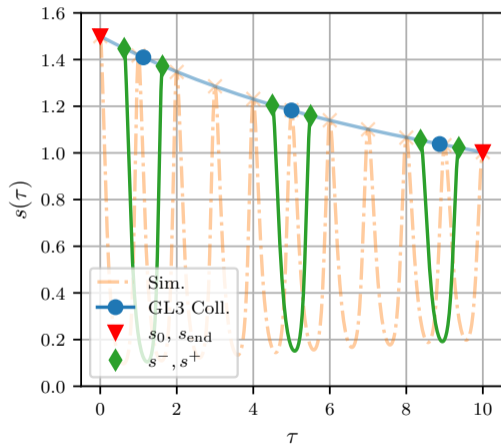
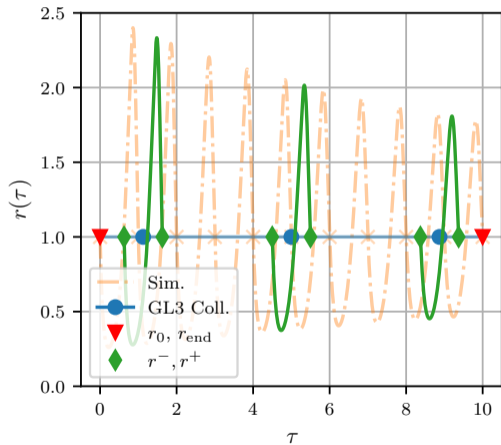
Example Integration: Predator-Prey IVP



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Control Parametrization





- ▶ Control matrix

$$U(\tau) = [u_0(\tau), u_1(\tau), \dots, u_{N_{\text{ctr}}-1}(\tau)] \in \mathbb{R}^{n_u \times N_{\text{ctr}}}$$



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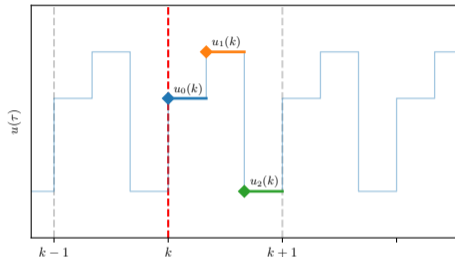
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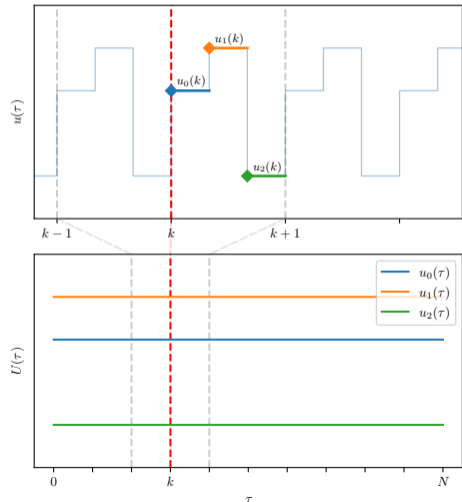
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$$U(\tau) = U_{\text{const}}$$



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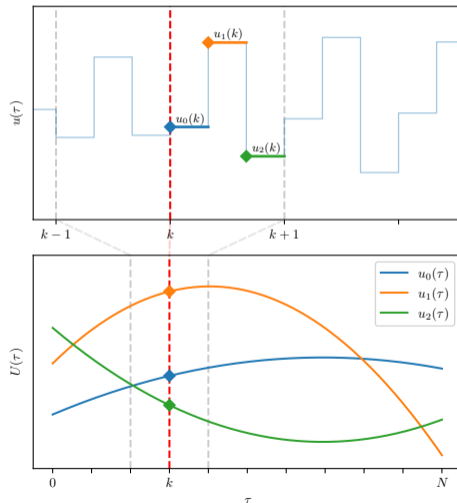
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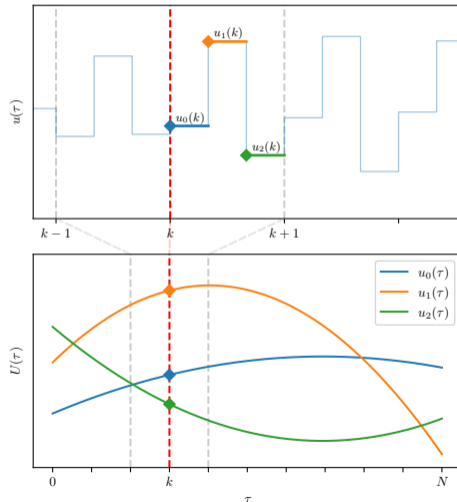
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- ▶ DAE:

$$\dot{x} = f_e(x(\tau), z(\tau), U(\tau))$$

$$0 = g_e(x(\tau), z(\tau), U(\tau))$$





Example: Predator Prey OCP

$$\begin{aligned} \min_{x_e(\cdot), z_e(\cdot), U(\cdot)} \quad & \|x_e(N) - \bar{x}_N\|_Q^2 \\ \text{s.t.} \quad & 0 = x_e(0) - \bar{x}_0, \\ & \dot{x}_e = f_e(x_e(\tau), z_e(\tau), U(\tau)), \quad \tau \in [0, N], \\ & 0 = g_e(x_e(\tau), z_e(\tau), U(\tau)), \quad \tau \in [0, N], \\ & 0 \leq h(x_e(\tau), z_e(\tau), U(\tau)), \quad \tau \in [0, N] \end{aligned}$$

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- Drive populations to target state $\bar{x}_N = [1, 0.7]^\top$.

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- ▶ Drive populations to target state $\bar{x}_N = [1, 0.7]^\top$.
- ▶ Discretize with 3 stage Gauss-Legendre collocation scheme over $N = 20$ periods.

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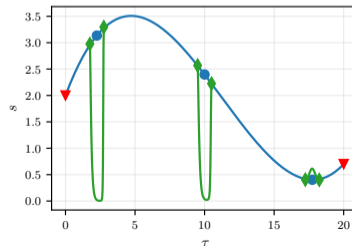
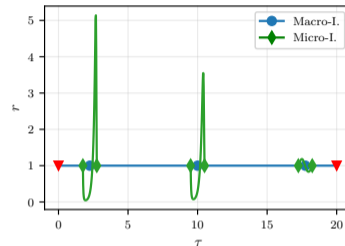
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- ▶ Drive populations to target state $\bar{x}_N = [1, 0.7]^\top$.
- ▶ Discretize with 3 stage Gauss-Legendre collocation scheme over $N = 20$ periods.
- ▶ The target state cannot be reached!

Example: Predator Prey OCP

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 & && 0 \leq h(x_e(\tau), z_e(\tau), U(\tau)), \quad \tau \in [0, N]
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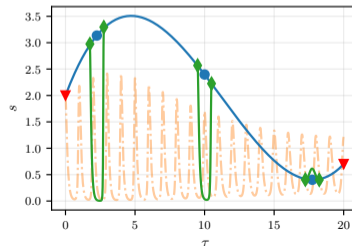
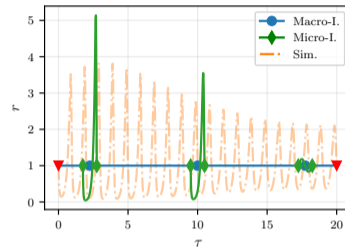
- ▶ Drive populations to target state $\bar{x}_N = [1, 0.7]^\top$.
- ▶ Discretize with 3 stage Gauss-Legendre collocation scheme over $N = 20$ periods.
- ▶ The target state cannot be reached!



Example: Predator Prey OCP

$$\begin{aligned}
 & \min_{x_e(\cdot), z_e(\cdot), U(\cdot)} && \|x_e(N) - \bar{x}_N\|_Q^2 \\
 & \text{s.t.} && 0 = x_e(0) - \bar{x}_0, \\
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Highly Oscillatory Trajectories



- ▶ Large approximation and integration error if the envelope is not smooth enough

Highly Oscillatory Trajectories

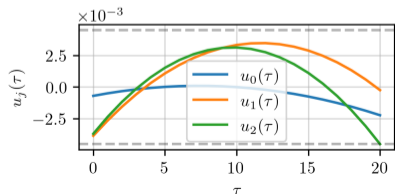
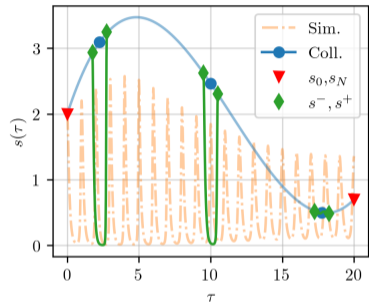


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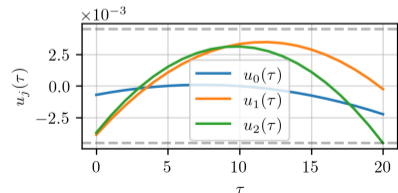
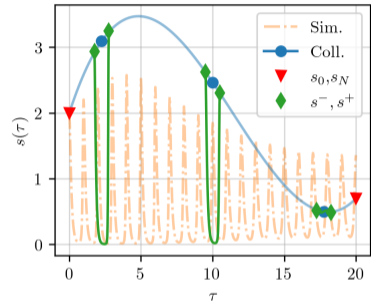


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- ▶ Ex: Regularization on 2nd envelope derivative

$$J = \|x_{\text{end}} - \bar{x}_{\text{end}}\|_Q^2 + \alpha \sum_{i=1}^d b_i w_i^2$$

where $w_i = \ddot{x}_e(\tau_i)$

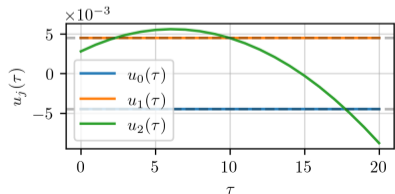
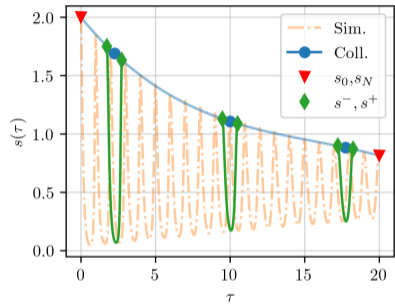


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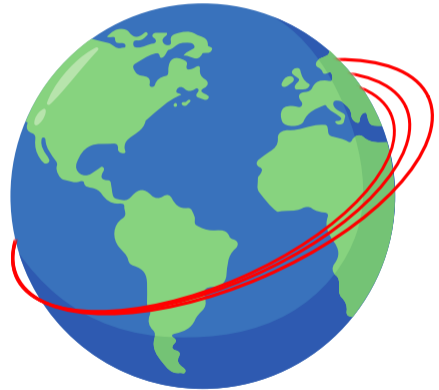
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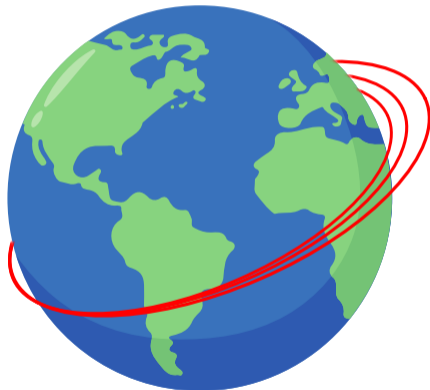
Example: Low Thrust Satellite



Graphics from Vecteezy.com

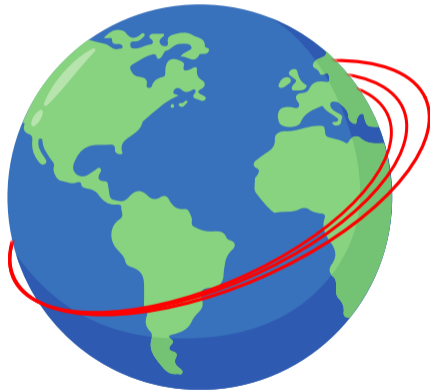
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- ▶ Move satellite from circular park orbit at 500 km to mission orbit at 1400 km, problem of [1].



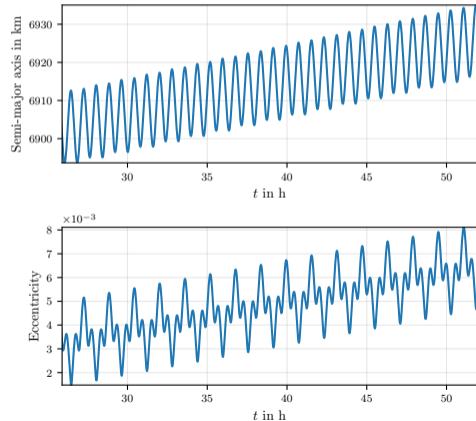
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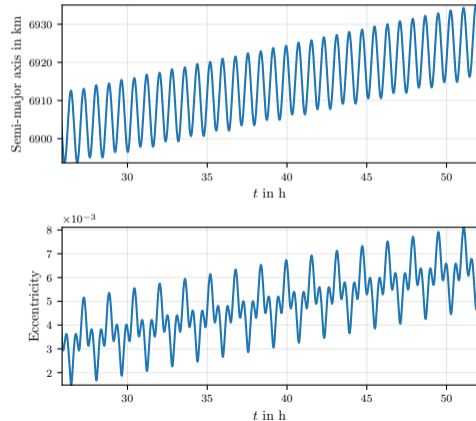
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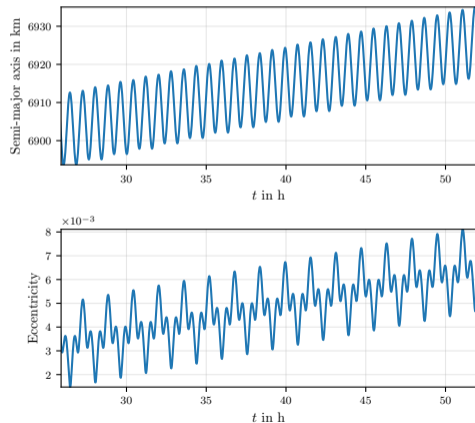
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- ▶ Here: Two stages with $N = 289$ orbits, we search for an optimal periodic control scheme with $N_{\text{ctrl}} = 30$ controls for each orbit in both stages



Example: Low Thrust Satellite



- ▶ State: Modified Equinoctial Elements

$$x = [p, f, g, h, k, L, m, t]^T \in \mathbb{R}^8$$



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$$\min_{\substack{x_e(\cdot), z_e(\cdot), \\ U_1, U_2}} m(2N)$$

s.t.

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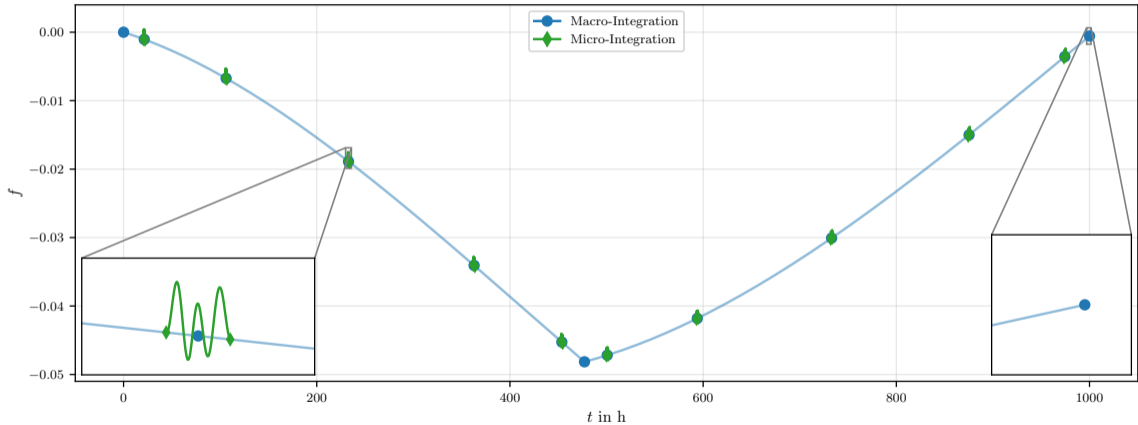
$$\dot{x} = f_e(x_e, z_e, U_2), \quad \tau \in [N, 2N],$$

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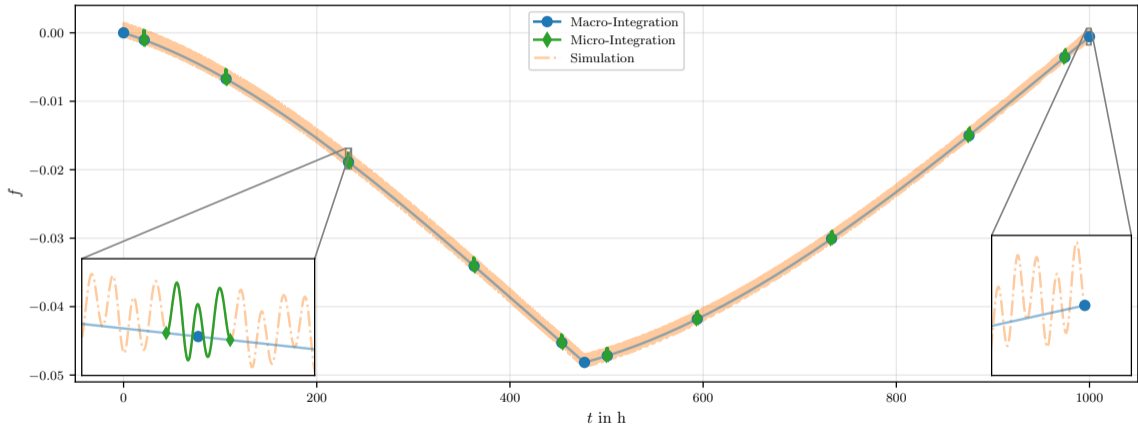
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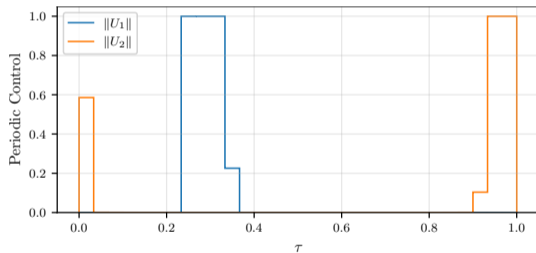
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Example: Low Thrust Satellite - Results

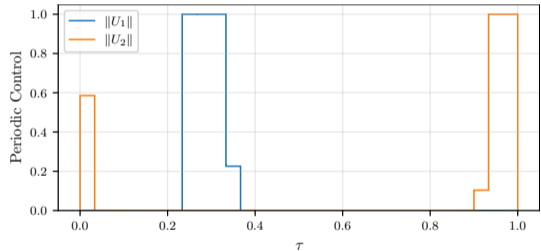


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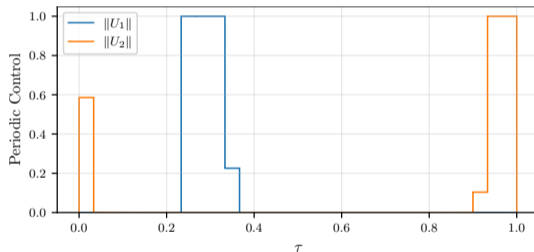
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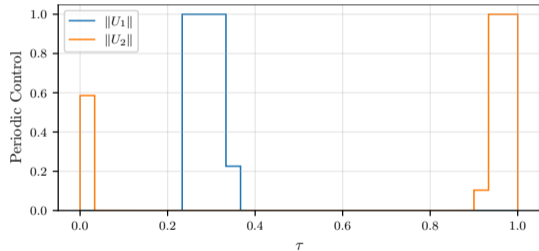


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- ▶ Optimal final mass $m = 155.48$ kg almost coincides with result of [1]



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Thank you for your attention!

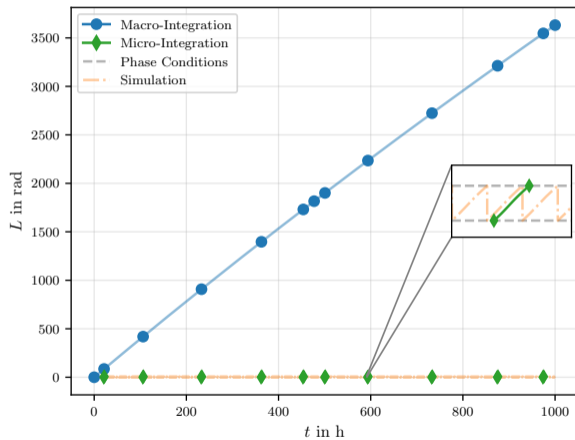
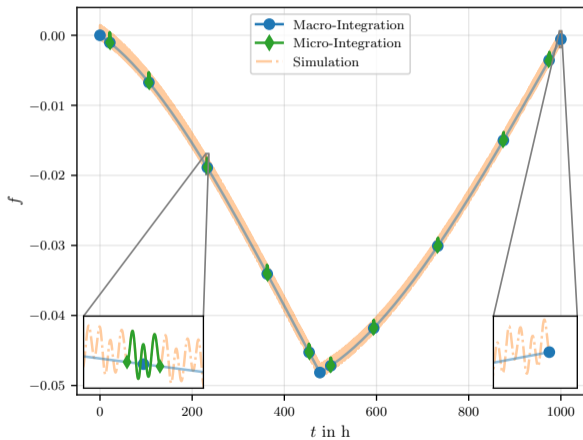


- [1] John T. Betts. “Very low-thrust trajectory optimization using a direct SQP method”. In: *Journal of Computational and Applied Mathematics* 120.1 (2000), pp. 27–40. ISSN: 0377-0427. DOI: [https://doi.org/10.1016/S0377-0427\(00\)00301-0](https://doi.org/10.1016/S0377-0427(00)00301-0). URL: <https://www.sciencedirect.com/science/article/pii/S0377042700003010>.
- [2] M. P. Calvo et al. “Numerical stroboscopic averaging for ODEs and DAEs”. In: *Applied Numerical Mathematics* 61 (Nov. 2010). DOI: 10.1016/j.apnum.2011.06.007.
- [3] M.P. Calvo. *Numerical stroboscopic averaging for ODEs and DAEs*. <https://www.fields.utoronto.ca/programs/scientific/11-12/SciCADE2011/presentations/Calvo.pdf>. [Online; accessed 1-November-2021]. 2011.
- [4] Otis Graf. “Multirevolution methods for orbit integration”. In: *Proceedings of the Conference on the Numerical Solution of Ordinary Differential Equations*. Ed. by Dale G. Bettis. Berlin, Heidelberg: Springer Berlin Heidelberg, 1974, pp. 471–490. ISBN: 978-3-540-37911-9.

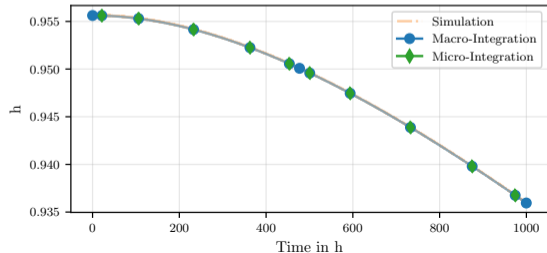
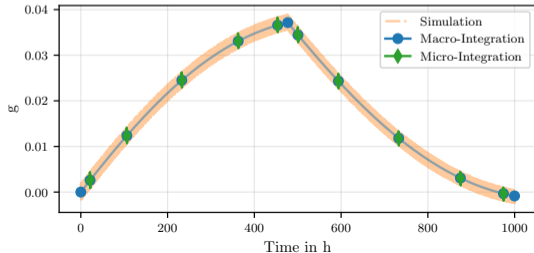
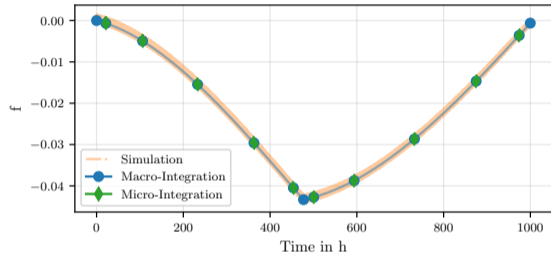
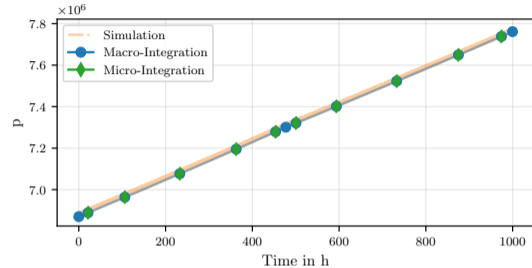


- [5] David Mace and L. H. Thomas. “An extrapolation formula for stepping the calculation of the orbit of an artificial satellite several revolutions time”. In: *The Astronomical Journal* 65 (1960), p. 300.
- [6] Linda R. Petzold. “An Efficient Numerical Method for Highly Oscillatory Ordinary Differential Equations”. In: *SIAM Journal on Numerical Analysis* 18.3 (1981), pp. 455–479. DOI: 10.1137/0718030. eprint: <https://doi.org/10.1137/0718030>. URL: <https://doi.org/10.1137/0718030>.

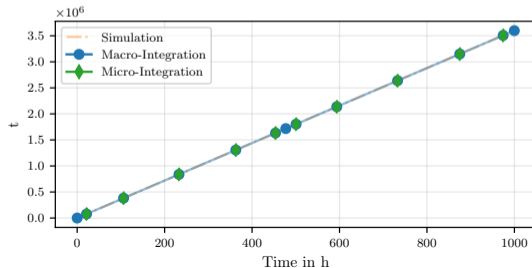
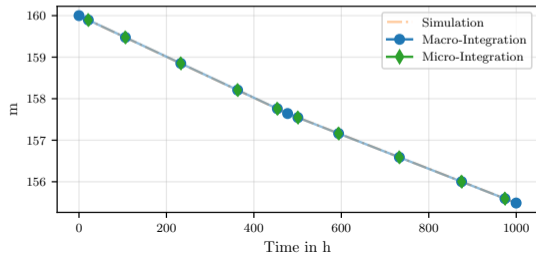
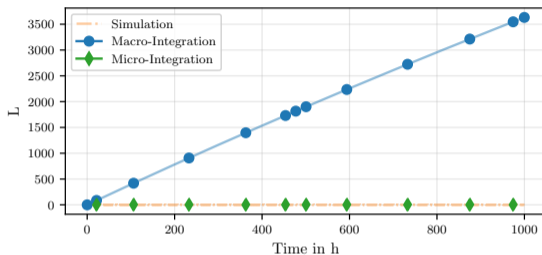
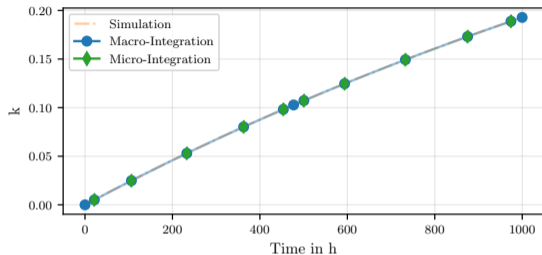
Example: Low Thrust Satellite - Solution



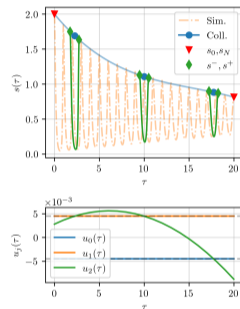
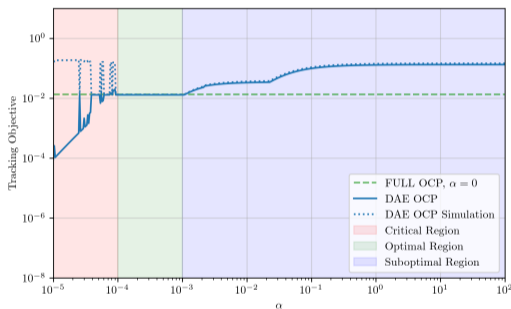
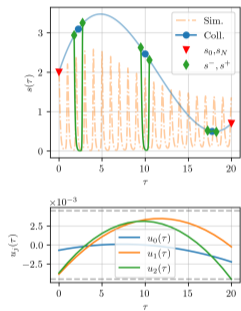
Example: Low Thrust Satellite - States 1



Example: Low Thrust Satellite - States 2



Regularization on Second State Derivative Solves the Problem



Left (red): too small regularization

Right (green): sufficiently high regularization

Note: "sweet spot" around $\alpha = 0.5 \cdot 10^{-3}$ where objective is unaffected by regularization

Macro-Integration of the DAE

- ▶ Integrate DAE from x_0 over N periods
- ▶ Single implicit d -stage Runge-Kutta step with butcher coefficients c, A, b :

$$0 = x_0 - \bar{x}_0$$

$$0 = x_N - x_0 + N \sum_{j=1}^d b_j v_j$$

$$0 = x_i - x_0 + N \sum_{j=1}^d a_{i,j} v_j \quad i = 1, \dots, d$$

$$0 = v_i - f_e(x_i, z_i) \quad i = 1, \dots, d$$

$$0 = g_e(x_i, z_i) \quad i = 1, \dots, d$$

Butcher tableau:

c_1	$a_{1,1}$	$a_{1,2}$	\dots	$a_{1,d}$
c_2	$a_{2,1}$	$a_{2,2}$	\dots	$a_{2,d}$
\vdots	\vdots	\vdots	\ddots	\vdots
c_d	$a_{d,1}$	$a_{d,2}$	\dots	$a_{d,d}$
	b_1	b_2	\dots	b_d

Forward Euler, $d = 1$:

0	0
	1

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$$\min_w \|x_N - \bar{x}_N\|_Q^2$$

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$$0 = x_i - x_0 + N \sum_{j=1}^d a_{i,j} v_j \quad i = 1, \dots, d,$$

$$0 = v_i - f_e(x_i, z_i, U_{\text{const}}) \quad i = 1, \dots, d,$$

$$0 = g_e(x_i, z_i, U_{\text{const}}) \quad i = 1, \dots, d,$$

$$0 \leq h(w)$$

► $w = (x_0, x_1, \dots, x_d, v_1, \dots, v_d, z_1, \dots, z_d, x_N, U_{\text{const}})$

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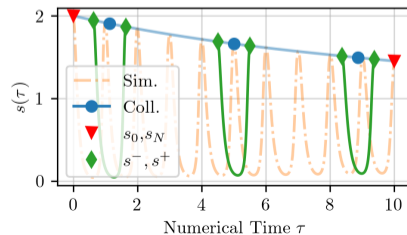
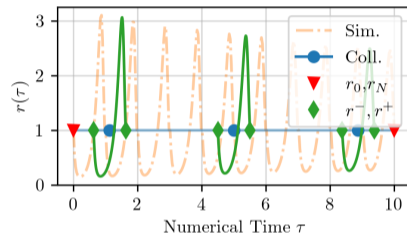
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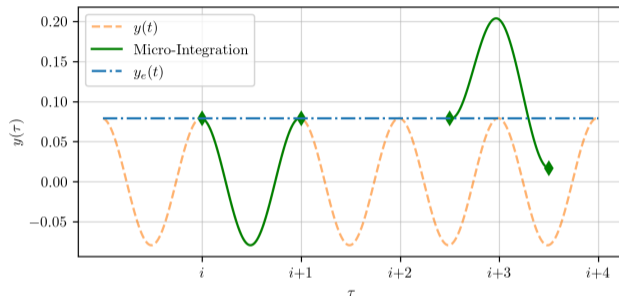
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$$0 \leq h(w)$$

► $U_{\text{const}}^* = [-2, 2, 2] \cdot 10^{-3}$





$$\dot{x} = \begin{bmatrix} \dot{y} \\ \dot{t} \end{bmatrix} = \begin{bmatrix} -y - \sin(\omega t) \\ 1 \end{bmatrix} \quad (3)$$

$$x(\tau) = \begin{bmatrix} x(\tau) \\ t(\tau) \end{bmatrix} = \begin{bmatrix} Ce^{-T\tau} - \frac{\sin(2\pi\tau)}{\omega^2+1} + \frac{\omega \cos(2\pi\tau)}{\omega^2+1} \\ T\tau \end{bmatrix} \quad (4)$$

Numerical Time

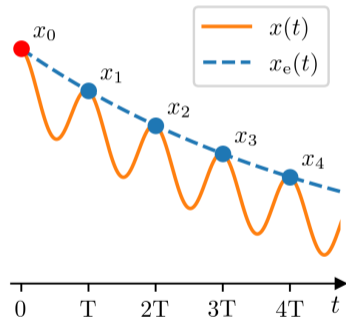
- ▶ Physical time: $t \in [0, t_f] \rightarrow$ Numerical time: $\tau \in [0, N]$
- ▶ Scaling with period duration T :

$$\tau(t) = \frac{t}{T} \qquad t(\tau) = \tau T$$

- ▶ Trajectories now in numerical time as $x(\tau)$, dynamics read

$$\frac{dx}{d\tau}(\tau) = T f(x(\tau))$$

- ▶ With numerical time, we can integrate over whole periods
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